

Assignment 4

1. Find the following for the wave $\psi = 5 \cos(2x + 3t)$
- wave vector
 - wavelength
 - frequency
 - period
 - phase velocity
 - amplitude

Wave vector $k = 2 \text{ meters}^{-1}$

Wavelength $\lambda = \frac{2\pi}{k} = \pi \text{ meters}$

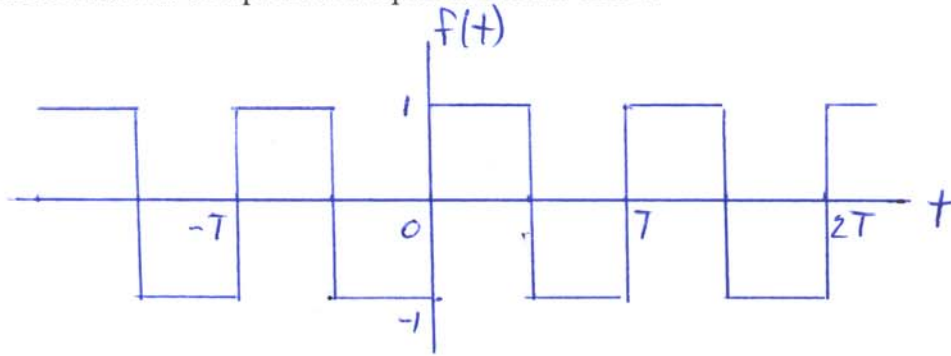
Frequency $\nu = \frac{3}{2\pi} \text{ sec}^{-1}$

Period $T = \frac{1}{\nu} = \frac{2\pi}{3} \text{ sec.}$

Phase Velocity $v_p = \frac{\omega}{k} = \frac{3}{2} \text{ meter/sec}$

Amplitude $A = 5 \text{ meters}$

2. Consider a series of square wave pulses shown below.



Fourier analysis says that this pulse can be expanded as:

$$f(t) = \frac{A_0}{2} + \sum_{n=1}^{\infty} (A_n \cos n\omega t + B_n \sin n\omega t) \quad \omega \equiv \frac{2\pi}{T}$$

- a) Show that $A_n = 4/n\pi$ where n is odd, $A_n = 0$ for n is even
 $B_n = 0$ and $A_0 = 0$

$$f(t) \text{ is odd} \Rightarrow A_0 + A_n = 0.$$

$$\therefore f(t) = \sum_{n=1}^{\infty} B_n \sin n\omega t$$

$$\int_0^T f(t) \sin m\omega t dt = \sum_{n=1}^{\infty} B_n \underbrace{\int_0^T \sin n\omega t \sin m\omega t dt}_{\equiv I} \quad (1)$$

Integral I

$$\begin{aligned} \text{Case 1: } n=m \quad I &= \int_0^T \sin^2 n\omega t dt \\ &= \int_0^T \frac{1 - \cos 2n\omega t}{2} dt \\ &= \left[\frac{t}{2} - \frac{\sin 2n\omega t}{4n\omega} \right]_0^T \\ &= \frac{T}{2} \end{aligned}$$

Case 2: $n \neq m$

$$I = \int_0^T \frac{\cos(n-m)\omega t - \cos(n+m)\omega t}{2} dt$$

$$= \left[\frac{\sin(n-m)\omega t}{2(n-m)\omega} - \frac{\sin(n+m)\omega t}{2(n+m)\omega} \right]_0^T$$

$$\therefore \int_0^T \sin n\omega t \sin m\omega t dt = \frac{T}{2} \delta_{nm}$$

$$\int_0^T f(t) \sin n\omega t dt = \int_0^{T/2} \sin n\omega t dt - \int_{T/2}^T \sin m\omega t dt$$

$$= \left[-\frac{\cos n\omega t}{n\omega} \right]_0^{T/2} + \left[\frac{\cos n\omega t}{n\omega} \right]_{T/2}^T$$

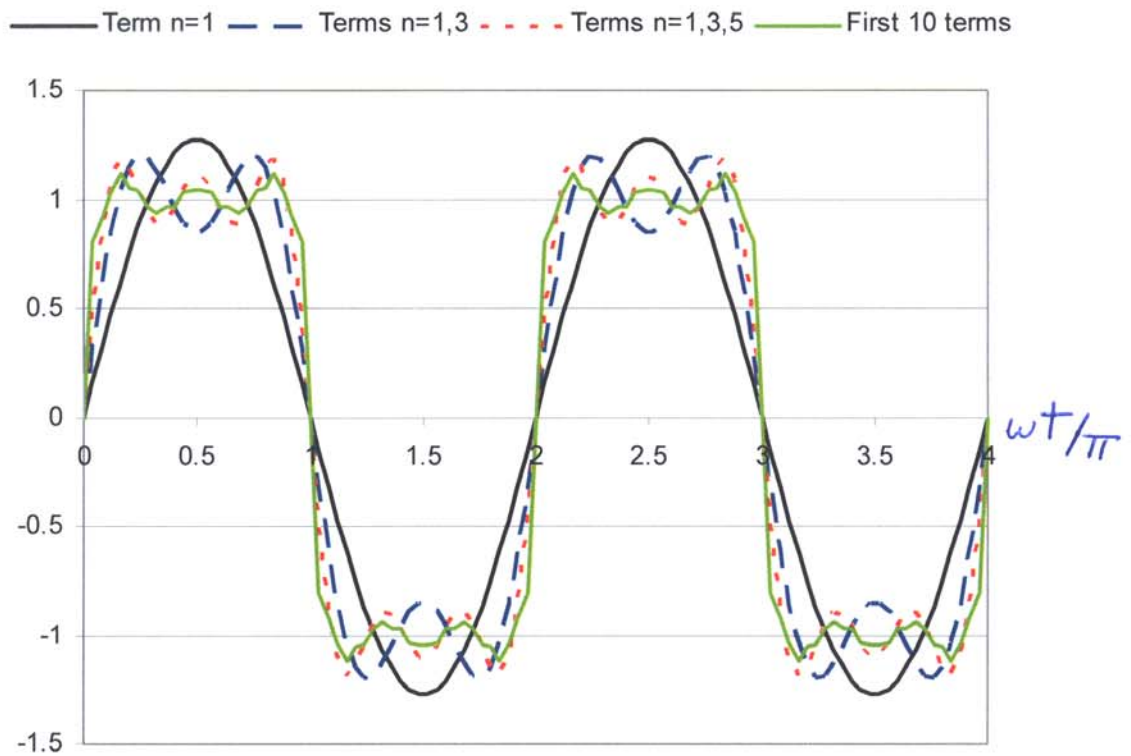
$$= \frac{1}{n\omega} \left(1 - \cos \frac{n\omega T}{2} \right) + \frac{1}{n\omega} \left(\cos n\omega T - \cos \frac{n\omega T}{2} \right)$$

$$= \begin{cases} 0 & \text{for } n \text{ even} \\ \frac{4}{n\omega} & \text{for } n \text{ odd.} \end{cases}$$

$$\therefore (i) \Rightarrow \frac{4}{n\omega} = B_n \frac{T}{2} \Rightarrow B_n = \begin{cases} 0 & \text{for } n \text{ even} \\ \frac{4}{n\pi} & \text{for } n \text{ odd.} \end{cases}$$

$$\therefore f(t) = \frac{4}{\pi} \sum_{n=1,3,5} \frac{\sin n\omega t}{n}$$

b) Plot the first term, first 2 terms and first 3 terms in the sum.



c) Hence, a pulse of light which can convey information, is composed of many frequencies. Estimate the range of frequencies $\Delta\nu$ required to make a one femtosecond laser pulse using the Heisenberg Uncertainty Principle $\Delta\nu \Delta t > 2\pi$.

$$\Delta t = 10^{-15} \text{ sec.}$$

$$\Rightarrow \Delta\nu \approx \frac{2\pi}{10^{-15}} \\ \approx 6 \times 10^{15} \text{ Hz.}$$

Hence, a femtosecond pulse does not have a well defined frequency!

3. Superposition Principle

- a) Show that if ψ_1 and ψ_2 are solutions of the 3 dimensional wave equation that their sum also is a solution.

$$\begin{aligned}\nabla^2 (\psi_1 + \psi_2) &= \nabla^2 \psi_1 + \nabla^2 \psi_2 \\ &= \frac{1}{v^2} \frac{d^2 \psi_1}{dt^2} + \frac{1}{v^2} \frac{d^2 \psi_2}{dt^2} \\ &= \frac{1}{v^2} \frac{d^2}{dt^2} (\psi_1 + \psi_2)\end{aligned}$$

$\therefore \psi_1 + \psi_2$ is also a solution.

- b) This may seem trivial but show that the superposition principle does not hold for the following differential equation.

$$\frac{d^2 \psi}{dx^2} = \psi^2$$

Let ψ_1 & ψ_2 be solns. i.e. $\frac{d^2 \psi_1}{dx^2} = \psi_1^2$

$$\frac{d^2 \psi_2}{dx^2} = \psi_2^2$$

$$\text{But } \frac{d^2}{dx^2} (\psi_1 + \psi_2) = \frac{d^2 \psi_1}{dx^2} + \frac{d^2 \psi_2}{dx^2}$$

$$= \psi_1^2 + \psi_2^2$$

$$\neq (\psi_1 + \psi_2)^2$$

$\therefore \psi_1 + \psi_2$ is not a solution.

4. Damped harmonic oscillator

$$m \frac{d^2 x}{dt^2} = -kx - \gamma \frac{dx}{dt}$$

- a) Consider a solution $x = A e^{\lambda x}$. Solve for λ . (Result will be complex) This approach is much simpler than using $x = A \cos \omega t + B \sin \omega t$.
- b) Write down the general solution for the case of weak damping $km \gg \gamma^2$.
- c) What is the solution for the case the mass is initially at rest at distance x_0 ? Plot this solution.

$$x = A e^{\lambda x} \Rightarrow m \lambda^2 = -k - \gamma \lambda$$

$$m \lambda^2 + \gamma \lambda + k = 0.$$

$$\lambda^2 + \frac{\gamma}{m} \lambda + \frac{k}{m} = 0$$

$$\lambda = \frac{-\gamma}{m} \pm \frac{\sqrt{\frac{\gamma^2}{m^2} - \frac{4k}{m}}}{2}$$

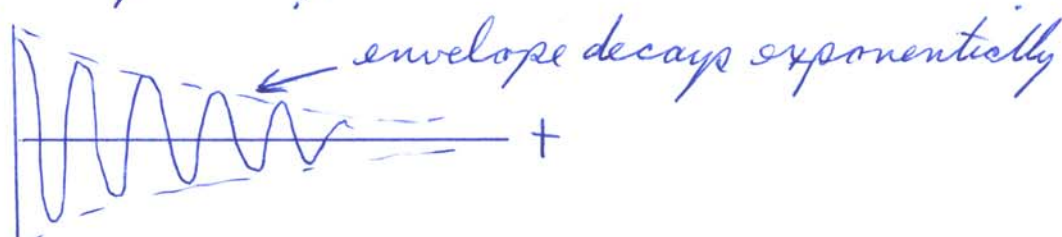
$$\lambda = \frac{-\gamma}{2m} \pm i \sqrt{\frac{k}{m} - \frac{\gamma^2}{4m^2}}$$

$$\approx \frac{-\gamma}{2m} \pm i \omega_0 \quad \omega_0 \equiv \sqrt{k/m}$$

$$\therefore x = e^{-\gamma t / 2m} \left[A e^{i \omega_0 t} + B e^{-i \omega_0 t} \right]$$

$$= e^{-\gamma t / 2m} \left[(A+B) \cos \omega_0 t + i(A-B) \sin \omega_0 t \right]$$

Taking real part yields: $x = x_0 e^{-\gamma t / 2m} \cos \omega_0 t$



5. Show that the group velocity v_g is related to the phase velocity v by the following equation. Note that for the case of normal dispersion $v_g < v$.

$$v_g = \frac{c}{n + \omega \frac{dn}{d\omega}}$$

Group Velocity $v_g \equiv \frac{d\omega}{dk}$

Phase Velocity $v_p = \frac{c}{n} = \frac{\omega}{k}$

$$\text{or } \omega = \frac{ck}{n}$$

$$\begin{aligned} \therefore \frac{d\omega}{dk} &= \frac{c}{n} - \frac{ck}{n^2} \frac{dn}{dk} \\ &= \frac{c}{n} - \frac{ck}{n^2} \frac{dn}{d\omega} \cdot \frac{d\omega}{dk} \end{aligned}$$

$$\frac{d\omega}{dk} \left(1 + \frac{ck}{n^2} \frac{dn}{d\omega} \right) = \frac{c}{n}$$

$$v_g = \frac{c/n}{1 + \frac{ck}{n^2} \frac{dn}{d\omega}}$$

$$\therefore v_g = \frac{c}{n + \omega \frac{dn}{d\omega}}$$