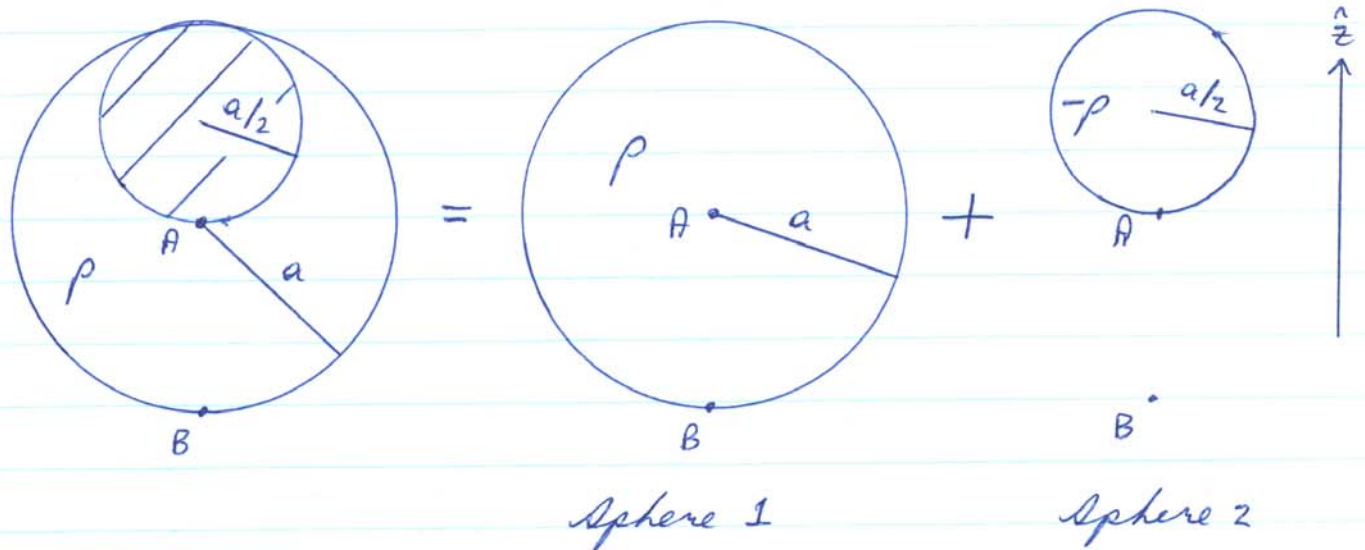


PHYS 2020 Assignment 3

1. Using the superposition principle the problem may be divided into two simpler problems.



Charge on sphere 1 $q_1 = \frac{4}{3} \pi a^3 \rho$

" " 2 $q_2 = -\frac{4}{3} \pi \left(\frac{a}{2}\right)^3 \rho$

Let $\vec{E}_1$ & $\vec{E}_2$ denote electric fields produced by q_1 & q_2 .

$$\therefore \vec{E}(A) = \vec{E}_1(A) + \vec{E}_2(A)$$

$$= 0 + \frac{q_2 (-\hat{z})}{(a/2)^2}$$

As expected $\vec{E}_2 \parallel \hat{z}$
since $q_2 < 0$.

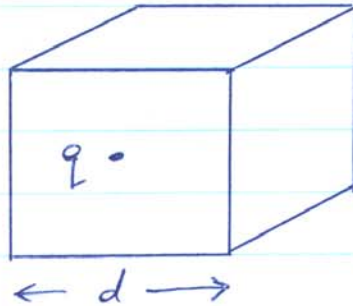
$$= \frac{2\pi}{3} a \rho \hat{z}$$

$$\vec{E}(B) = \vec{E}_1(B) + \vec{E}_2(B)$$

$$= \frac{q_1}{a^2} (-\hat{z}) + \frac{q_2}{(a + a/2)^2} (-\hat{z})$$

$$= \frac{-34}{27} \pi a \rho \hat{z}.$$

2a)



$$\int \vec{E} \cdot d\vec{a} = 4\pi q$$

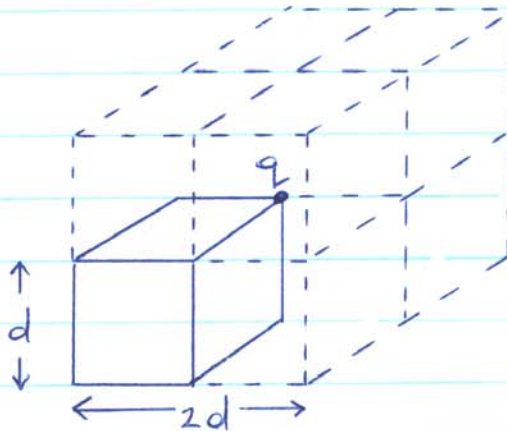
all surfaces
of cube

Since charge is at cube center, an equal amount of electric field flux passes through each of 6 sides.

$$\therefore \int \vec{E} \cdot d\vec{a} = \frac{4\pi q}{6} = \frac{2\pi q}{3}$$

one face
of cube

b)



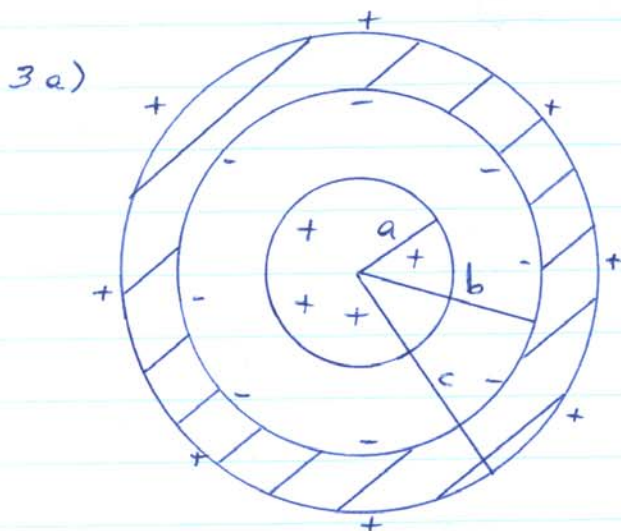
Since $\vec{E}$ extends radially outward from q , it does not cross the 3 planes that touch q at one corner. i.e. $\vec{E} \cdot d\vec{a} = 0 \Rightarrow \int \vec{E} \cdot d\vec{a} = 0$
side touching q

An equal amount of flux passes through each of the remaining 3 sides. (left, front & bottom)
To find this flux we construct a cube of length $2d$ centered about q .

$$\int_{\substack{\text{all surfaces} \\ \text{of length } 2d}} \vec{E} \cdot d\vec{a} = 4\pi q$$

$$4 \times 6 \int_{\substack{\text{a } d \times d \text{ surface} \\ \text{of cube of length } 2d}} \vec{E} \cdot d\vec{a} = 4\pi q$$

$$\therefore \int \vec{E} \cdot d\vec{a} = \frac{\pi q}{6} \text{ of remaining 3 surfaces.}$$



$$r < a \quad \rho = \frac{\rho_0 r}{a} \text{ insulator}$$

$$a < r < b \quad \text{vacuum}$$

$$b < r < c \quad \text{conductor}$$

$$c < r \quad \text{vacuum}$$

$\vec{E} = 0$ inside conductor. Hence a negative surface charge must exist at $r = b$ to counteract the charged insulator. Since the conductor is neutral an equal positive charge must reside on the outer conductor surface at $r = c$.

i) $\vec{E} = E(r) \hat{r}$ due to spherical symmetry

$$\int_{\text{surface of sphere of radius } r} \vec{E} \cdot d\vec{a} = 4\pi \int_{\text{volume of sphere of radius } r} \rho \, dV$$

$$r < a \quad E(r) 4\pi r^2 = 4\pi \int_0^r \frac{\rho_0 r}{a} 4\pi r^2 \, dr$$

$$= 4\pi \frac{4\pi \rho_0}{a} \int_0^r r^3 \, dr$$

$$= 4\pi \frac{4\pi \rho_0}{a} \frac{r^4}{4}$$

$$\therefore \vec{E}(r) = \frac{\pi \rho_0}{a} r^2 \hat{r}$$

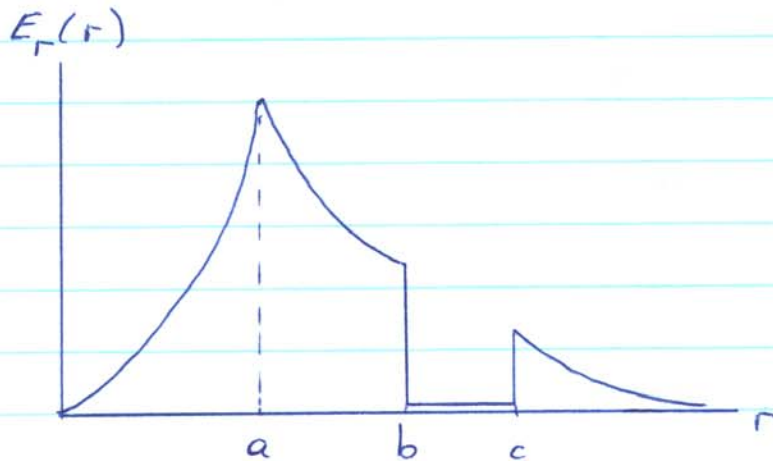
Note that charge on insulator $q = \pi \rho_0 a^3$.

$$a < r < b \quad E(r) \quad 4\pi r^2 = 4\pi q$$

$$\vec{E}(r) = \frac{q}{r^2} \hat{r}$$

$$b < r < c \quad \vec{E} = 0 \text{ inside conductor}$$

$$r > c \quad \vec{E} = \frac{q}{r^2} \hat{r}$$



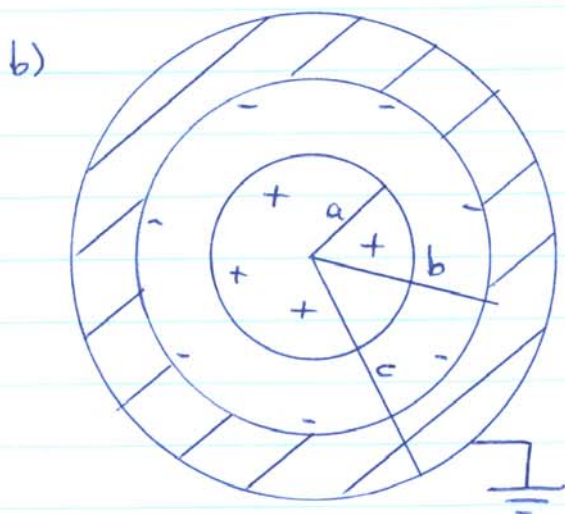
ii) Electric Field Outside Conductor $\vec{E} = 4\pi \epsilon \hat{n}$

where $\hat{n}$ = direction perpendicularly outward from surface.

$$\begin{aligned} \sigma(r=b) &= \frac{\vec{E}(b)}{4\pi} \cdot (-\hat{r}) \quad (\text{Outward direction from conductor at } r=b \text{ is } -\hat{r}) \\ &= \frac{1}{4\pi} \frac{q}{b^2} \hat{r} \cdot (-\hat{r}) \\ &= -\frac{q}{4\pi b^2} \\ &= -\frac{\rho_0 a^3}{4b^2} \end{aligned}$$

$$\begin{aligned}
 \sigma(r=c) &= \frac{\vec{E}(c) \cdot \hat{r}}{4\pi} \\
 &= \frac{1}{4\pi} \frac{q}{c^2} \hat{r} \cdot \hat{r} \\
 &= \frac{\rho_0 a^3}{4c^2}
 \end{aligned}$$

Note that total amount of charge on conductor
 $= \sigma(r=b) \cdot 4\pi b^2 + \sigma(r=c) \cdot 4\pi c^2$
 $= 0.$



Since $\vec{E} = 0$ inside a conductor there still is the negative charge at the inner conductor surface. However since the conductor is connected to ground the outer positive charge flows away into ground.

$$\therefore \sigma(r=b) = -\frac{\rho_0 a^3}{4b^2} \quad (\text{same as before})$$

$$\sigma(r=c) = 0.$$

Note that $\sigma(r=b) \cdot 4\pi b^2 = -\pi \rho_0 a^3$
 $= -q \quad (q = \text{charge of insulator})$

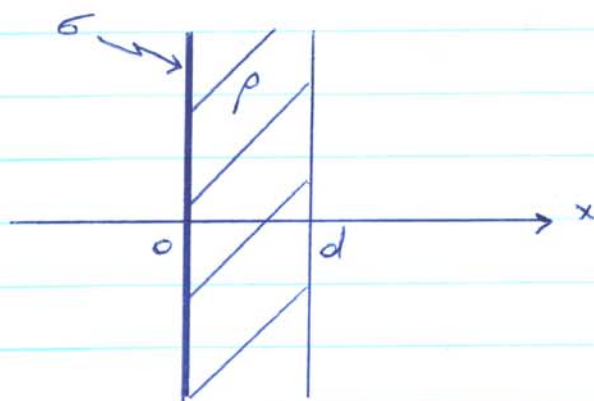
$\therefore$ when $r > c$ using Gauss law we get:

$$E(r) 4\pi r^2 = 4\pi [q + (-q)]$$

$$\vec{E}(r) = 0.$$

For $r < c$, electric field is the same as in part a.

4.



By symmetry $\vec{E} \parallel \hat{x}$.

In class we showed that an infinite uniformly charged plane generates a uniform field. Hence $\vec{E}$ has the same magnitude but opposite direction on either side of the charge strip.

Case 1: $x < 0$ or $x > d$.

Consider a cylindrical pillbox straddling the strip



$$\int_{\text{cylinder surface}} \vec{E} \cdot d\vec{a} = 4\pi \int_{\text{cylinder volume}} \rho dV.$$

$$EA + EA = 4\pi (\sigma + \rho d) A$$

$$E = 2\pi (\sigma + \rho d)$$

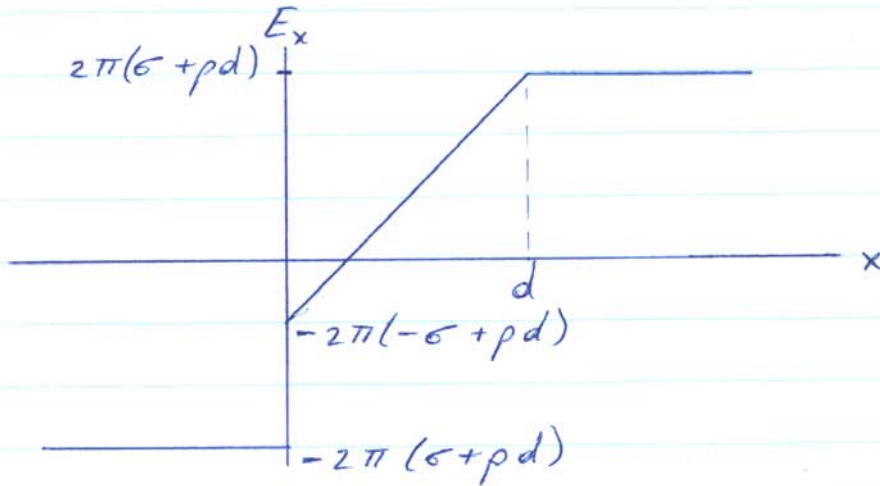
This result expresses that uniformly charged infinite planes attract or repel charges depending on which side of the plane the charge is.

$$\therefore E(x) = -2\pi (\text{charge per cm}^2 \text{ to right of } x) + 2\pi (\text{charge per cm}^2 \text{ to left of } x)$$

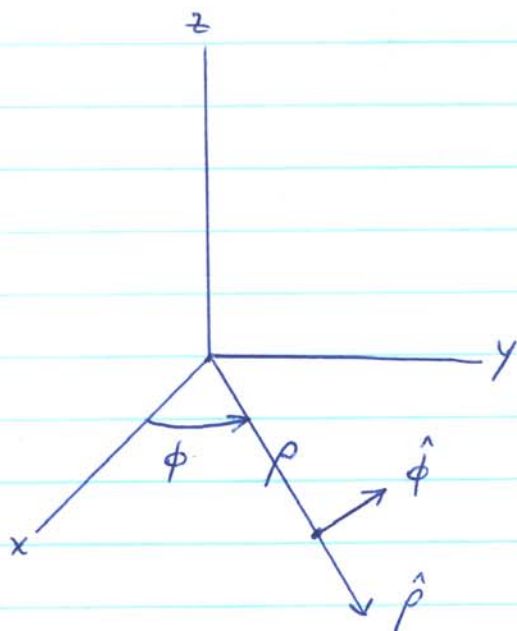
Case 2: $0 < x < d$.

$$E(x) = -2\pi \rho(d-x) + 2\pi(\sigma + \rho x)$$

$$= 2\pi(\sigma - \rho d + 2x).$$



5 i)



$$z = z$$

$$y = \rho \sin \phi$$

$$x = \rho \cos \phi$$

$$ii) \quad \hat{z} = \hat{z} \quad (1)$$

$$\hat{\rho} = (\cos \phi, \sin \phi, 0) = \hat{x} \cos \phi + \hat{y} \sin \phi \quad (2)$$

$$\hat{\phi} = (-\sin \phi, \cos \phi, 0) = -\hat{x} \sin \phi + \hat{y} \cos \phi. \quad (3)$$

Solving (2) + (3) for $\hat{x}$ + $\hat{y}$ we get:

$$\hat{x} = \hat{\rho} \cos \phi - \hat{\phi} \sin \phi \quad (4)$$

$$\hat{y} = \hat{\rho} \sin \phi + \hat{\phi} \cos \phi \quad (5)$$

$$iii) \quad E_x = \vec{E} \cdot \hat{x}$$

$$= \vec{E} \cdot (\hat{\rho} \cos \phi - \hat{\phi} \sin \phi)$$

$$E_x = E_\rho \cos \phi - E_\phi \sin \phi.$$

$$E_y = \vec{E} \cdot \hat{y}$$

$$= \vec{E} \cdot (\hat{\rho} \sin \phi + \hat{\phi} \cos \phi)$$

$$E_y = E_\rho \sin \phi + E_\phi \cos \phi$$

$$iv) \frac{\partial E_y}{\partial y} = \frac{\partial E_y}{\partial \rho} \frac{\partial \rho}{\partial y} + \frac{\partial E_y}{\partial \phi} \frac{\partial \phi}{\partial y} + \frac{\partial E_y}{\partial z} \frac{\partial z}{\partial y}$$

$$\frac{\partial E_z}{\partial z} = \frac{\partial E_z}{\partial \rho} \frac{\partial \rho}{\partial z} + \frac{\partial E_z}{\partial \phi} \frac{\partial \phi}{\partial z} + \frac{\partial E_z}{\partial z} \frac{\partial z}{\partial z}$$

$$v) \rho = \sqrt{x^2 + y^2}$$

$$\frac{\partial \rho}{\partial x} = \frac{1}{2} (x^2 + y^2)^{-1/2} \cdot 2x = \frac{x}{\rho} = \cos \phi$$

$$\frac{\partial \rho}{\partial y} = \frac{1}{2} (x^2 + y^2)^{-1/2} \cdot 2y = \frac{y}{\rho} = \sin \phi$$

$$\frac{\partial \rho}{\partial z} = 0$$

$$\tan \phi = \frac{y}{x}$$

$$\sec^2 \phi \frac{\partial \phi}{\partial x} = -\frac{y}{x^2}$$

$$= -\frac{\sin \phi}{\rho \cos^2 \phi}$$

$$\frac{\partial \phi}{\partial x} = -\frac{\sin \phi}{\rho}$$

$$\sec^2 \phi \frac{\partial \phi}{\partial y} = \frac{1}{x}$$

$$= \frac{1}{\rho \cos \phi}$$

$$\frac{\partial \phi}{\partial y} = \frac{\cos \phi}{\rho}$$

Obviously $\frac{\partial \phi}{\partial z} = 0.$

$$\frac{\partial z}{\partial x} = \frac{\partial z}{\partial y} = 0 \quad \frac{\partial z}{\partial z} = 1.$$

vi) $\nabla \cdot \vec{E} = \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z}$

$$\frac{\partial E_x}{\partial x} = \frac{\partial E_x}{\partial \rho} \cos \phi + \frac{\partial E_x}{\partial \phi} \left(\frac{-\sin \phi}{\rho} \right) + 0 \quad \text{from (iv) + (v)}$$

$$= \frac{\partial}{\partial \rho} (E_\rho \cos \phi - E_\phi \sin \phi) \cos \phi$$

$$+ \frac{\partial}{\partial \phi} (E_\rho \cos \phi - E_\phi \sin \phi) \left(\frac{-\sin \phi}{\rho} \right) \quad \text{using (iii)}$$

$$\frac{\partial E_x}{\partial x} = \left(\frac{\partial E_\rho}{\partial \rho} \cos \phi - \frac{\partial E_\phi}{\partial \rho} \sin \phi \right) \cos \phi$$

$$+ \left(\frac{\partial E_\rho}{\partial \phi} \cos \phi - E_\rho \sin \phi - \frac{\partial E_\phi}{\partial \phi} \sin \phi - E_\phi \cos \phi \right) \left(\frac{-\sin \phi}{\rho} \right)$$

$$\frac{\partial E_y}{\partial y} = \frac{\partial E_y}{\partial \rho} \sin \phi + \frac{\partial E_y}{\partial \phi} \frac{\cos \phi}{\rho} + 0 \quad \text{from (iv) + (v)}.$$

$$= \frac{\partial}{\partial \rho} (E_\rho \sin \phi + E_\phi \cos \phi) \sin \phi$$

$$+ \frac{\partial}{\partial \phi} (E_\rho \sin \phi + E_\phi \cos \phi) \frac{\cos \phi}{\rho}$$

$$\frac{\partial E_y}{\partial y} = \left(\frac{\partial E_r}{\partial \rho} \sin \phi + \frac{\partial E_\phi}{\partial \rho} \cos \phi \right) \sin \phi$$

$$+ \left(\frac{\partial E_r}{\partial \phi} \sin \phi + E_r \cos \phi + \frac{\partial E_\phi}{\partial \phi} \cos \phi - E_\phi \sin \phi \right) \frac{\cos \phi}{\rho}$$

$$\therefore \nabla \cdot \vec{E} = \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z}$$

$$= \frac{\partial E_r}{\partial \rho} (\cos^2 \phi + \sin^2 \phi)$$

$$+ \frac{\partial E_\phi}{\partial \rho} (-\sin \phi \cos \phi + \cos \phi \sin \phi)$$

$$+ \frac{\partial E_r}{\partial \phi} \left(-\cos \phi \frac{\sin \phi}{\rho} + \sin \phi \frac{\cos \phi}{\rho} \right)$$

$$+ \frac{\partial E_\phi}{\partial \phi} \left(\frac{\sin^2 \phi}{\rho} + \frac{\cos^2 \phi}{\rho} \right)$$

$$+ E_r \left(\frac{\sin^2 \phi}{\rho} + \frac{\cos^2 \phi}{\rho} \right)$$

$$+ E_\phi \left(\cos \phi \frac{\sin \phi}{\rho} - \sin \phi \frac{\cos \phi}{\rho} \right)$$

$$+ \frac{\partial E_z}{\partial z}$$

$$\nabla \cdot \vec{E} = \frac{\partial E_r}{\partial \rho} + \frac{1}{\rho} \frac{\partial E_\phi}{\partial \phi} + \frac{E_r}{\rho} + \frac{\partial E_z}{\partial z}$$

$$\text{or } \nabla \cdot \vec{E} = \frac{1}{\rho} \frac{\partial (\rho E_r)}{\partial \rho} + \frac{1}{\rho} \frac{\partial E_\phi}{\partial \phi} + \frac{\partial E_z}{\partial z}$$